



# Metastable states during dissociation of carbon dioxide hydrates below 273 K

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## ABSTRACT

In this study, the dissociation of isolated carbon dioxide hydrate particles of sizes in the range 0.25–2.5 mm was investigated. It was found that below the ice melting point, the hydrates dissociated into supercooled water (metastable liquid) and gas. The formation of the liquid phase during CO<sub>2</sub> hydrate dissociation was visually observed, and the pressures of the hydrate dissociation into supercooled water and gas were measured in the temperature range 249–273 K. These pressures agreed well with the calculated data for the supercooled water–hydrate–gas metastable equilibrium (Istomin et al., 2006). In the *P*–*T* area on the phase diagram bounded by the ice–hydrate–gas equilibrium curve and the supercooled water–hydrate–gas metastable equilibrium curve, hydrates could exist for a long time because the metastable phase and their stability are not connected to the self-preservation effect. The growth of the metastable CO<sub>2</sub> hydrate film on the surface of supercooled water droplets formed during the hydrate dissociation was observed at pressure above the three-phase supercooled water–hydrate–gas metastable equilibrium pressure but still below the three-phase ice–hydrate–gas equilibrium pressure. It was found that the growth rate of the metastable CO<sub>2</sub> hydrate film was higher by a factor of 25 and 50 than that for methane hydrate and propane hydrate, respectively.

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## 1. Introduction

The unexpected stability of clathrate hydrates of some gases at temperatures below the ice melting point outside their stable region, known as the self-preservation effect (Yakushev and Istomin, 1992) or the anomalous preservation regime (Stern et al., 2001), has also been observed for CO<sub>2</sub> hydrates (Chuvilin et al., 2007; Circone et al., 2003; Falenty and Kuhs, 2009; Kuhs et al., 2004; Takeya and Ripmeester, 2008). The particular interest in the stability and dissociation of CO<sub>2</sub> hydrates below 273 K is due to the development of a conceptual idea for CO<sub>2</sub> sequestration as a hydrate in permafrost to reduce the emission of carbon dioxide into the atmosphere and abate global warming (Duchkov et al., 2009; Goel, 2006).

Because the anomalous stability of hydrates outside their stable region was observed only below the ice melting point, it was proposed that the effect was caused by the formation of an ice coating on the surface of hydrate particles (Davidson et al., 1986; Yakushev and Istomin, 1992). Recently, we detected (Melnikov et al., 2010) the existence of metastable methane and propane hydrates below the three-phase ice–hydrate–gas equilibrium pressure, and the stability of the hydrates was not related with ice

formation on their surface. In the pressure–temperature area between the ice–hydrate–gas equilibrium curve and the supercooled water–hydrate–gas metastable equilibrium curve (Fig. 1) hydrate particles that were 0.15–2.5 mm in diameter existed for a prolonged time without visually observed evidence of hydrate dissociation. The hydrates were dissociated into supercooled water and gas when crossing the supercooled water–hydrate–gas metastable equilibrium curve once the pressure was decreased or the temperature was increased. The formation of supercooled water on dissociation of methane hydrate was detected in the range 253–273 K (Melnikov et al., 2009) and 257–273 K on dissociation of propane hydrate (Mel'nikov et al., 2007).

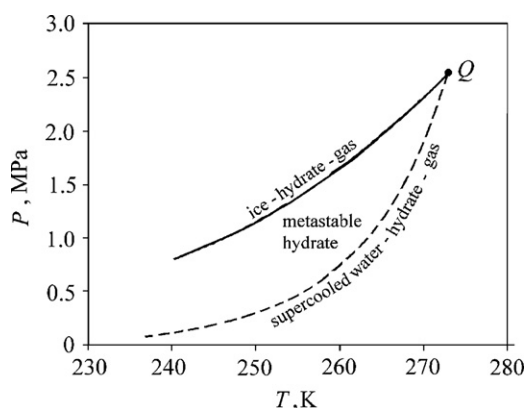
In this communication, the results of a study on behavior of CO<sub>2</sub> hydrate particles outside the hydrate stable region below the ice melting point are presented.

## 2. Experimental section

CO<sub>2</sub> gas (99.9 mol%) and distilled water were used to synthesize hydrate samples. A detailed description of the experimental setup used in this work and the hydrate synthesis procedure are given elsewhere (Melnikov et al., 2009; Mel'nikov et al., 2007). The main element of the experimental setup was a stainless steel cylindrical reactor with a volume of 100 ccm. For visual observation, the lateral surface of the reactor was equipped with two quartz windows

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**Fig. 1.** Phase diagram for the  $\text{H}_2\text{O}-\text{CH}_4$  system. Solid line is the ice-hydrate-gas equilibrium curve and dashed line is the supercooled water-hydrate-gas metastable equilibrium curve. Q is the quadruple point, where ice+water+hydrate+gas coexist (modified from Melnikov et al. (2010)).

located opposite to one another at the half-height of the reactor. Hydrates were formed from water sprayed on a transparent Plexiglas plate. The plate was previously cooled to 253–255 K to form frozen droplets of 0.25–2.5 mm diameter and about 0.3 mm thickness. The plate with frozen droplets was mounted vertically inside the reactor between the windows. The pressure in the reactor was measured by a transducer with an accuracy of 0.01 MPa. Two copper-constantan thermocouples were used to measure the temperature inside the reactor at the top and bottom, and the average temperature was calculated. The temperature inside the reactor was maintained with an accuracy of  $\pm 0.1$  K. The reactor with frozen water droplets was evacuated at 255 K and charged with  $\text{CO}_2$  gas to a pressure between the equilibrium pressure of  $\text{CO}_2$  hydrate and the dew point of carbon dioxide at given temperature.

To increase the rate of hydrate formation, the reactor with gas and water droplets was slowly heated from 255 to 274 K. The melting/freezing procedure (heating of the reactor to 274 K and subsequent cooling to 255 K) was repeated a few times to provide additional conversion of water into hydrates. The same procedure was used by us before synthesizing methane and propane hydrates from frozen water droplets (Melnikov et al., 2009; Mel'nikov et al., 2007).

For optical observation, a cathetometer was used. The cathetometer telescope was fitted with a digital camera, and the picture was recorded in DVD in real time.

To observe the dissociation of hydrates, the pressure in the reactor was slowly reduced at given temperature. The dissociation of the hydrates was judged from the visually observed collapse of the rough surface of the hydrates, the appearance of smooth islands of the liquid phase on hydrate particles, and the evolution of gas bubbles from the liquid. The pressure at which the first changes in hydrate particles were visually observed was taken as the hydrate dissociation pressure  $P_d$  at given temperature. After the onset of hydrate dissociation, the subsequent observations were carried out at constant pressure  $P_d$ . It should be noted that the influence of hydrate particle size on dissociation pressure  $P_d$  was not detected. All the particles were dissociated at the same pressure  $P_d$ .

### 3. Results and discussion

#### 3.1. Supercooled water- $\text{CO}_2$ hydrates-gas metastable equilibrium

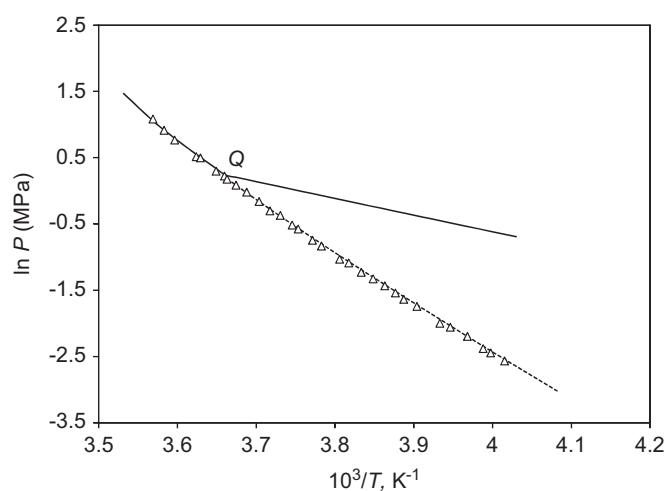
The experimental data on  $\text{CO}_2$  hydrate dissociation pressure  $P_d$  obtained in this study are summarized in Table 1 and plotted in Fig. 2.

Above the quadruple point temperature,  $P_d$  is in good agreement with  $\text{CO}_2$  hydrate equilibrium pressure calculated by CSMGem program (Sloan and Koh, 2007).

**Table 1**

Experimental dissociation pressure  $P_d$  for  $\text{CO}_2$  hydrates formed from water droplets.

$T$ (K)	$P$ (MPa)
$T > 273$ K	
280.2	2.95
279.1	2.49
278.1	2.15
276.0	1.67
274.1	1.34
273.3	1.25
$T < 273$ K	
272.2	1.09
271.2	0.98
270.0	0.85
269.0	0.74
268.1	0.69
267.0	0.60
266.4	0.56
265.2	0.48
264.0	0.43
262.8	0.36
262.0	0.34
260.9	0.29
259.9	0.27
258.9	0.24
258.0	0.21
257.3	0.19
256.2	0.17
254.3	0.14
253.4	0.13
252.0	0.11
250.8	0.09
249.1	0.08



**Fig. 2.** Dissociation pressure of  $\text{CO}_2$  hydrates formed from water droplets (symbols) and the equilibrium pressure of hydrate formation for bulk  $\text{CO}_2$  hydrates (—) calculated by the CSMGem program (Sloan and Koh, 2007). Dashed line is the calculated extension of the water-hydrate-methane gas equilibrium curve to the metastable area of supercooled water (Istomin et al., 2006).

Below the quadruple point temperature, the  $P_d$  values fall on the supercooled water- $\text{CO}_2$  hydrate-gas metastable equilibrium curve calculated by Istomin et al. (2006). These can be expressed in the form of regression equations as follows:

$$\ln(P_d/P_0) = 28.61 - 7774.7/T, \quad r^2 = 0.9982$$

$$73 \text{ K} < T < 249 \text{ K}$$

(1)

where  $P_0 = 1$  MPa and  $r^2$  is the regression coefficient.

Fig. 3 shows the successive shots of changes observed visually during the dissociation of CO<sub>2</sub> hydrates at 266.4 K. The first changes in hydrate particles appeared at 0.56 MPa (Fig. 3b). Ice is the stable phase at these conditions, but we observed that the hydrates dissociated into supercooled (metastable) water and gas. The supercooled water lifetime in the metastable state from the beginning of hydrate dissociation until crystallization of supercooled water depended on hydrate dissociation temperature. At 270 K, this stage lasted for tens of hours, but it was a few seconds at 250 K. The absence of ice in the hydrate samples is a necessary constraint to detect supercooled water during hydrate dissociation. Supercooled water cannot coexist with ice, and it freezes instantaneously in the presence of ice. We used isolated small water droplets to synthesize the hydrate samples free from ice. Even though some of hydrate particles formed will contain the ice, they will not influence the behavior of other hydrate particles during their dissociation. It is very difficult to achieve full conversion of water into gas hydrate for bulk samples, even if powdered ice is used. We suppose that presence of ice in hydrate samples can explain that Falenty and Kuhs (2009) in their neutron diffraction experiments had no evidence for any substantial amount of water in the liquid state during the dissociation of bulk hydrate below the ice melting point.

We observed the formation of supercooled water during CO<sub>2</sub> hydrate dissociation up to 249 K. Below this temperature and below the CO<sub>2</sub> hydrate equilibrium pressure no visual changes in the hydrate samples were observed. If the hydrate samples, cooled below 249 K, were heated at a pressure below the hydrate equilibrium pressure, liquid water was observed only near the ice melting point.

### 3.2. Stability of CO<sub>2</sub> hydrates below the three-phase ice–hydrate–gas equilibrium pressure

We observed the existence of CO<sub>2</sub> hydrates for a long time without visible evidence of the hydrate dissociation in the area bounded by the ice–hydrate–gas equilibrium curve and the supercooled water–hydrate–gas metastable equilibrium curve. Fig. 4 shows the stable CO<sub>2</sub> hydrates (Fig. 4a) and the hydrates, which were stored for 90 h at 264 K and 0.6 MPa (Fig. 4b). The

CO<sub>2</sub> hydrate equilibrium pressure at 264 K is 0.78 MPa, but no visual evidence of hydrate dissociation was detected at 264 K and 0.6 MPa during the entire duration of hydrate observation (90 h). When the pressure in the reactor was decreased, we observed the appearance of supercooled water immediately following the intersection of the supercooled water–hydrate–gas metastable equilibrium curve (Fig. 4c). Based on this, we can conclude that the CO<sub>2</sub> hydrate samples did not contain ice, and the hydrate stability in the area bounded by the ice–hydrate–gas equilibrium curve and the curve of the supercooled water–hydrate–gas metastable equilibrium was not connected with the self-preservation effect. The hydrate observed in our experiments in the area bounded by the ice–hydrate–gas equilibrium curve and the curve of the supercooled water–hydrate–gas metastable equilibrium is an example of the metastable phase, the stability of which is caused by the kinetic difficulty in the direct transformation of the hydrate crystal structure into ice. The stability of the self-preserved hydrate is caused by the slow rate of gas diffusion through the ice layer on the surface of hydrate particles (Takeya et al., 2001).

### 3.3. Hydrate growth below the three-phase ice–hydrate–gas equilibrium pressure

Fig. 5 shows the lateral growth of the film of metastable CO<sub>2</sub> hydrates at the surface of a supercooled water droplet at 265.2 K and 0.57 MPa. Because of the stochastic nature of hydrate nucleation, hydrates are not formed at the same time on the droplets of supercooled water. In this experiment, after the completion of the hydrate dissociation into supercooled water and gas (hydrates were no longer detected in the water droplets, Fig. 5b), the pressure  $P_d=0.48$  MPa was kept in the reactor for the next 10 min and then it was increased to 0.56 MPa, but it still remained below the ice–hydrate–gas equilibrium pressure at 265.2 K, 0.82 MPa. About 1 min after the pressure was increased, the onset of growth of the solid film on the surface of the supercooled water droplet was observed (Fig. 5c). We checked that the solid film in Fig. 5c–f was hydrate by observing that it melted if the pressure was again lowered below  $P_d$ .

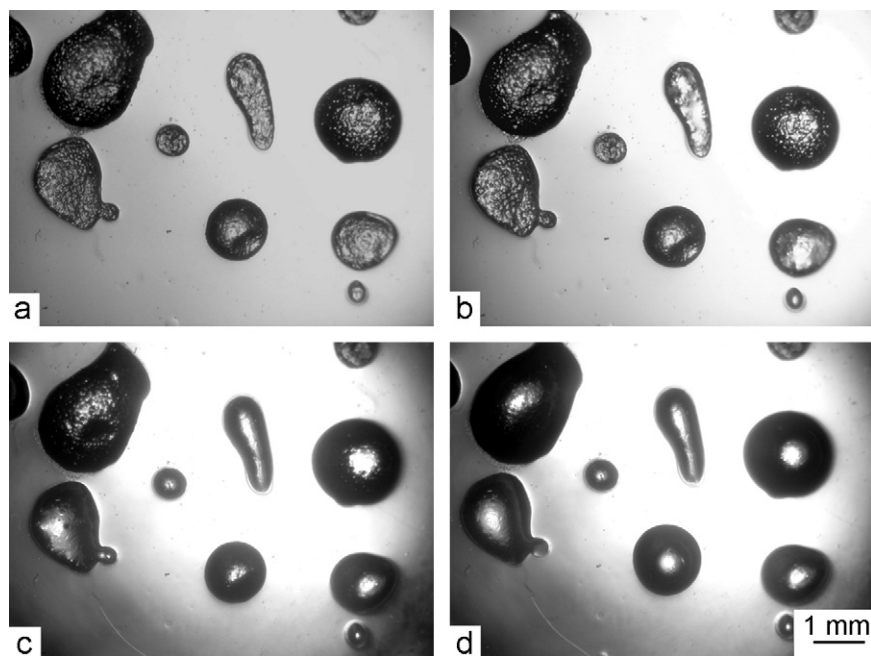
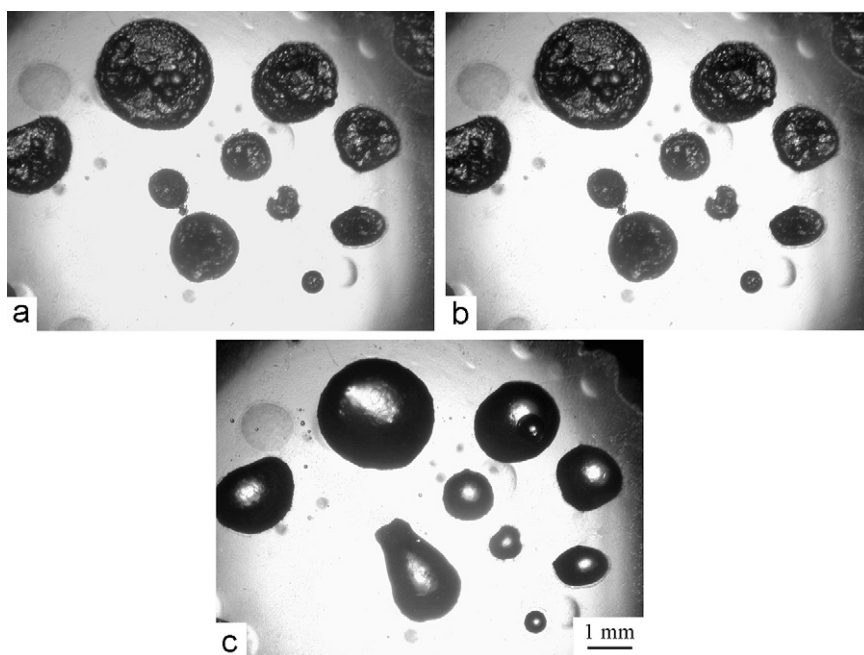
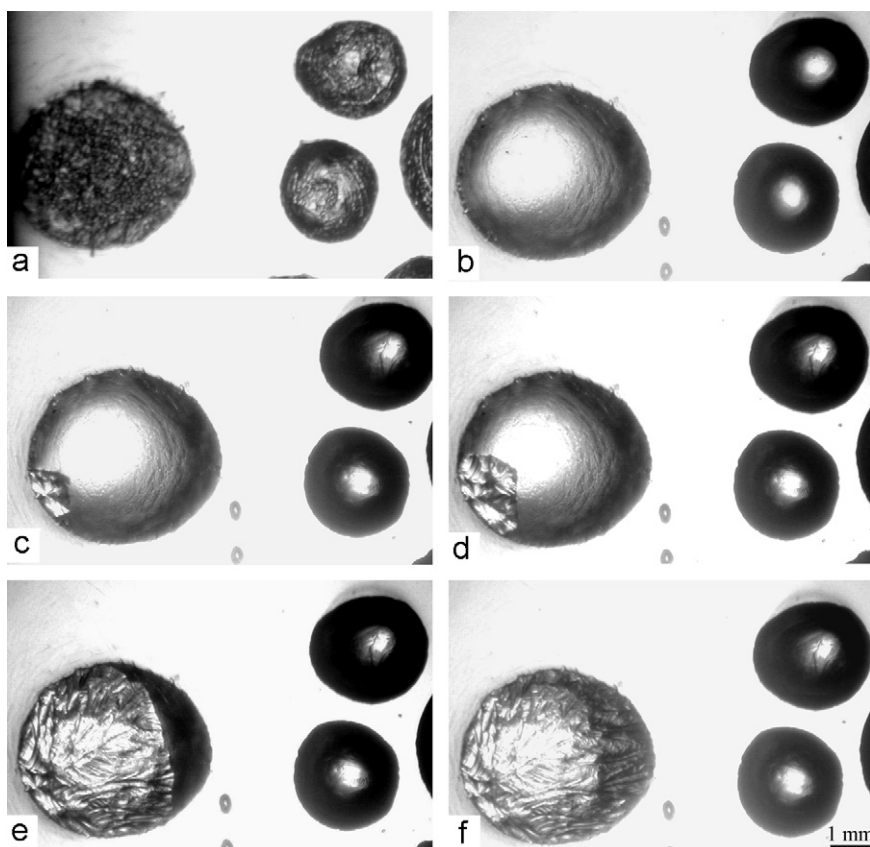


Fig. 3. Formation of supercooled water during CO<sub>2</sub> hydrate dissociation at 266.4 K. (a) At 1.8 MPa (before pressure decrease); (b) onset of hydrate dissociation at 0.56 MPa, time=0; (c) at 0.56 MPa, time=4 min; and (d) at 0.56 MPa, time=14 min.



**Fig. 4.** (a) Stable CO<sub>2</sub> hydrates at 2 MPa and 264 K; (b) metastable CO<sub>2</sub> hydrates 90 h after the pressure was decreased to 0.6 MPa at 264 K (the CO<sub>2</sub> hydrate equilibrium pressure is 0.78 MPa); and (c) dissociation of CO<sub>2</sub> hydrates into supercooled water and gas at 264 K and 0.43 MPa.



**Fig. 5.** (a and b) Dissociation of CO<sub>2</sub> hydrates into supercooled water and gas and (c–f) growth of metastable hydrates on the surface of supercooled water at 265.2 K. (a) Onset of hydrate dissociation at 0.48 MPa; (b) 15 min after the onset of hydrate dissociation,  $P=0.48$  MPa. Hydrates are no longer detected in the water droplets (the completion of the hydrate dissociation); (c–f) growth of metastable hydrates at 0.56 MPa. The pressure of the ice–hydrate–gas equilibrium at 265.2 K is 0.82 MPa.

With the use of video, the lateral growth rates of the CO<sub>2</sub> hydrate film on the surface of supercooled water droplets were measured. The procedure used was described in detail by Freer et al. (2001).

The lateral growth rates of the CO<sub>2</sub> hydrate film on the surface of supercooled water droplets as a function of supercooling  $\Delta T$  at different system temperatures are shown in Fig. 6. In these



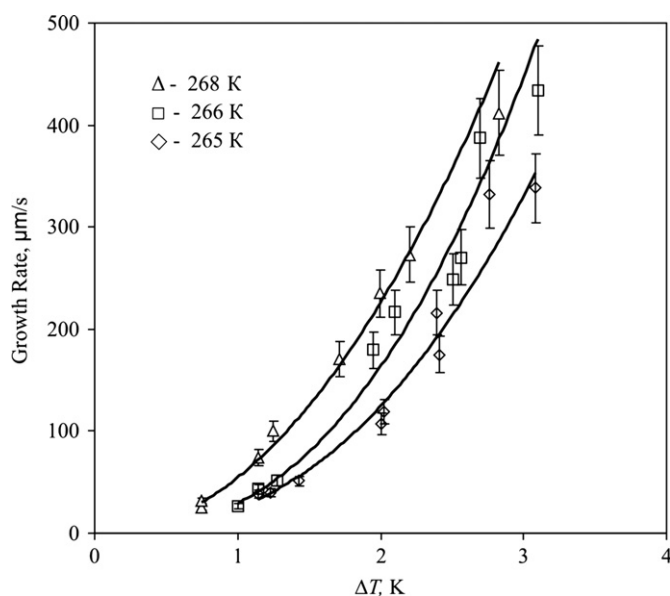


Fig. 6. Lateral growth rate of the film of metastable  $\text{CO}_2$  hydrate on the surface of supercooled water droplets formed during the hydrate dissociation.

experiments, water with a previous hydrate history (water obtained by melting  $\text{CO}_2$  hydrate at a pressure slightly below the supercooled water–hydrate–gas metastable equilibrium pressure) is used. Here  $\Delta T = T_{\text{meq}} - T$ , where  $T_{\text{meq}}$  is the temperature of the supercooled water–hydrate–gas metastable equilibrium at the system pressure  $P$ , and  $T$  is the system temperature. The comparison of the growth rate of the  $\text{CO}_2$  hydrate film on the surface of supercooled water droplets with the data on film growth for metastable methane and propane hydrates (Melnikov et al., 2010) demonstrates that the film growth rate for  $\text{CO}_2$  hydrate is about 25 times higher than that for methane hydrate and is about 50 times higher than that for propane hydrate. This difference is probably due to the higher solubility of  $\text{CO}_2$  gas into water in comparison with the solubility of methane and propane.

#### 4. Conclusions

The formation of liquid supercooled water as a transition state between  $\text{CO}_2$  hydrates and ice was visually observed during  $\text{CO}_2$  hydrate dissociation between 273 and 249 K. The pressure of the supercooled water–hydrate–gas metastable equilibrium was measured at different temperatures. Evidence for the existence of the metastable  $\text{CO}_2$  hydrates was obtained, and the stability of the hydrates was not connected with the ice formation on the hydrate surface in the area bounded by the ice–hydrate–gas equilibrium curve and the supercooled water–hydrate–gas metastable equilibrium curve. The lateral growth rates of the metastable  $\text{CO}_2$  hydrate film

on the surface of supercooled water were measured. It was found that the  $\text{CO}_2$  hydrate film growth rate was higher by a factor of 25 and 50 than those for methane hydrate and propane hydrate, respectively.

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